Rare, Exotic, and Invisible Higgs Decays at CMS



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h(125)



- A comprehensive program is underway to characterize the Higgs boson discovered in 2012
 - Mass and width
 - Spin-parity properties
 - Rates and kinematics of production and decay
- The Higgs boson provides a unique and wide ranging probe of Beyond the Standard Model (BSM) physics



• BSM Higgs boson decays and O(5%) deviations in properties may be the only observable sign of new physics at the LHC.

Rare, Exotic, and Invisible Higgs Decays

BSM Higgs Decays



Branching ratio to BSM decays is only constrained to <34% at 95% CL



Places to look:



Rare, Exotic, and Invisible Higgs Decays



Invisible

Compact Muon Solenoid





Rare, Exotic, and Invisible Higgs Decays

Higgs→x̄x (invisible)



 Higgs decays to undetected dark matter are inferred from the momentum imbalance of detected final state particles

$$\overrightarrow{\text{MET}} = \vec{E}_{\text{T}}^{\text{miss}} = -\sum_{\text{detected}} \vec{p}_{\text{T}}$$

• Multiple Higgs production channels



Higgs→xx̄ (invisible)

- Understanding MET distributions is crucial
- MET arises in standard model processes from neutrinos
 - e.g. $Z(v\bar{v})$ +jets
 - Modeling is sensitive to high order corrections
- MET also arises from momentum mis-measurement
- Extensive use of data control regions



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Data-Pred.) Data / Pred

Higgs→xx̄ (invisible) Results

 Observed (expected) 95% CL upper limits on invisible branching ratio from 13 TeV searches

	Data Sample	Observed BR	Expected BR
VBF	2.3 fb ⁻¹	< 69%	< 62%
Z(II)H	36 fb ⁻¹	< 40%	< 42%
Mono-jet	36 fb ⁻¹	<74%	< 57%
V(had)H aka Mono-V	36 fb ⁻¹	< 49%	< 45%

HIG-16-009, EXO-16-052, EXO-16-048

 The searches are statistically independent and are combined to improve sensitivity

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Higgs Portal Model



- Interpret in context of Higgs portal model
 - arXiv:0809.2745, arXiv:1212.2131, arXiv:1112.3299, arXiv:1205.3169
- Dark matter interacts with nuclei via Higgs boson exchange







pp→H*→HH



• The Higgs self coupling is determined by its potential



In the standard model,

 $V = \frac{m_h^2}{2}h^2 + \lambda_3 v h^3 + \frac{\lambda_4}{4}h^4$ where $\lambda_3 = \lambda_4 = m_h^2/(2v^2)$

- We can test λ_3 by studying di-Higgs production
 - Rare in standard model (cross section = 33.5 fb)
 - Constrain the coupling modifiers in these interfering diagrams:





The Hig LianTao Wang h^4 h O ,2Figure 8: Question of the nature of the electroweak phase transition. We can Understanding this physics is also directly relevant to one of the most fundamental questions we can ask about any symmetry breaking phenomenon, - Rare^Vin¹ which is what is the order of the associated phase transition. How can we experimentally decide whether the electroweak phase transition in the early Constra MS: universe was second order or first order? This quese ing shut and Tabyliu's talk ous next step following the Higgs discovery: having understood what breaks g Tuesday electroweak symmetry, we must now under take an experimental program to n Wednesday, August 13, 14 Tuesday, January 29, 15 probe how electroweak symmetry is restored at high energies. A first-order phase transition is also strongly motivated by the possibility of electroweak baryogenesis [18]. While the origin of the baryon asymmetry is one of the most fascinating questions in physics, it is frustratingly straightforward to build models for baryogenesis at ultra-high energy scales, with no direct experimental consequences. However, we aren't forced to defer this physics to the deep ultraviolet: as is well known the dynamics of electroweak - g symmetry breaking itself provides all the engredients needed for baryogeneh sis. At temperatures far above the weak scale, where electroweak symmetry

$pp \rightarrow H^* \rightarrow HH$

- Multiple final states
 - H(bb)H(bb)
 - H(bb)H(ττ)
 - H(bb)H(VV), leptonic V decays
 - H(bb)H(yy)
- Analysis sensitivity depends on branching ratios, backgrounds, and experimental resolutions
- The H(bb)H(yy) channel has the smallest branching ratio, but is the most sensitive to standard model di-Higgs production
 - $\sigma/\sigma_{SM} < 19.2$ at 95% CL

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$H(bb)H(\gamma\gamma)$



- Excellent di-photon mass resolution
- Fit data in 2D M($\gamma\gamma$)-M(bb) plane



$H(bb)H(\gamma\gamma)$ Results





Rare, Exotic, and Invisible Higgs Decays





H→4µ







Rare, E

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Summary



- A comprehensive program is underway to characterize the 125 GeV Higgs boson
- So far, searches for rare, exotic, and invisible Higgs decays show compatibility with standard model
- Recent results on
 - Higgs decays to undetected particles
 - Higgs self coupling
 - Higgs decays to light scalars
 - Lepton flavor violating decays
- Constraints are placed on a wide range of beyond the standard model physics models



Backup

Rare, Exotic, and Invisible Higgs Decays

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h(125) Properties (examples)





h(125) Properties (examples)





Rare, Exotic, and Invisible Higgs Decays

Higgs→xī (invisible)

- Understanding MET
 distributions is crucial
 - MET arises in standard model processes from neutrinos
 - e.g. Z(vv)+jets
 - Modeling is sensitive to high order corrections
 - MET also arises from momentum mis-measurement
 - Extensive use of data control regions
- Multivariate techniques
 - Boosted Decision Tree including MET, lepton kinematics, etc.
 - Improve sensitivity by ~10%

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Z(II)H(inv) Boosted Decision Tree



- $|m_{\ell\ell} m_Z|$ (dilepton mass);
- $p_{\rm T}^{\ell 1}$ (leading lepton transverse momentum);
- $p_{\rm T}^{\ell 2}$ (subleading lepton transverse momentum);
- $p_{\rm T}^{\ell\ell}$ (dilepton transverse momentum);
- $|\eta^{\ell 1}|$ (leading lepton pseudorapidity);
- $|\eta^{\ell 2}|$ (subleading lepton pseudorapidity);
- $E_{\rm T}^{\rm miss}$ (missing transverse energy);
- $m_T(p_T^{\ell 1}, E_T^{\text{miss}})$ (leading lepton transverse mass);
- $m_T(p_T^{\ell 2}, E_T^{\text{miss}})$ (subleading lepton transverse mass);
- $\Delta \phi(\vec{p_T}^{\ell\ell}, \vec{p_T}^{\text{miss}})$ (azimuthal separation between dilepton and missing energy);
- $\Delta R_{\ell\ell}$ (separation between leptons); and
- $|\cos \theta_{\ell 1}^{CS}|$ (cosine of Collins–Soper angle for leading lepton).







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H→4µ Dark SUSY Limits





Figure 6: 95% CL upper limits (black solid curves) from this search on $\sigma(pp \rightarrow h \rightarrow 2\gamma_D + X) \mathcal{B}(h \rightarrow 2\gamma_D + X)$ (with $m_{n_1} = 10 \text{ GeV}$, $m_{n_D} = 1 \text{ GeV}$) in the plane of two of the parameters (ε and m_{γ_D}) for the dark SUSY scenarios, along with constraints from other experiments. The colored contours represent different values of $\mathcal{B}(h \rightarrow 2\gamma_D + X)$ in the range 1–40%.

Image Credits



- Higgs potential: <u>https://inspirehep.net/record/1252561/plots</u>
- SUSY: <u>https://www.sciencenews.org/article/</u> <u>supersymmetry%E2%80%99s-absence-lhc-puzzles-</u> <u>physicists</u>
- Dark Matter: <u>https://map.gsfc.nasa.gov/universe/</u> <u>uni_matter.html</u>
- Extra Dimensions: <u>https://www.physics.uci.edu/~tanedo/</u> <u>docs.html</u>